

MASS EXCHANGE BETWEEN SPHERICAL BODIES  
AND A FLUID STREAM

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A solution of the problem of determining the mass-exchange coefficient between the surface of a sphere and a liquid flowing over the sphere at low Reynolds number (second approximation) is presented.

In 1952, the author and V. G. Levich independently obtained the following equation [1-3]:

$$\frac{\text{Nu}}{\sqrt[3]{\text{Pr}}} = 0.99 \sqrt[3]{\text{Re}}, \quad (1)$$

which is valid for  $\text{Re} < 1$ . The author used a diffusion integral relation governing the thickness of the diffusion layer:

$$\frac{\partial}{\partial \theta} \int_0^{\delta} (c - c_1) v_{\theta} dy + \frac{\text{ctg} \theta}{a} \int_0^{\delta} (c - c_1) v_{\theta} dy = -D \left( \frac{\partial c}{\partial y} \right)_{y=0}, \quad (2)$$

and the velocity distribution was determined from the known Stokes solution for the stream function

$$\psi = -\frac{U}{4} \sin^2 \theta \left( 2r^2 - 3ar + \frac{a^3}{r} \right).$$

Recently, Van Dyke [4] obtained a second approximation to describe the fluid flow near a sphere

$$\psi = -\frac{U}{4} \sin^2 \theta (r - a)^2 \left[ \left( 1 + \frac{3}{16} \text{Re} \right) \left( 2 + \frac{a}{r} \right) + \frac{3}{16} \text{Re} \left( 2 + \frac{a}{r} + \frac{a^2}{r^2} \right) \cos \theta \right]. \quad (3)$$

The possibility was therefore disclosed for obtaining the second approximation in the solution of the mass-exchange problem.

It follows from (3) that for  $\text{Re} \leq 16$  the tangential velocity near a sphere  $v_{\theta}$  is positive on the whole surface of the sphere. If  $\text{Re} > 16$ , this velocity takes on negative values at the rear part of the sphere for

$$\cos \theta_1 = -\left( \frac{3}{4} + \frac{4}{\text{Re}} \right) = -m. \quad (4)$$

In the first case we can follow the growth of the diffusion layer on the whole surface ( $0 \leq \theta \leq \pi$ ), on the major portion of the surface ( $0 \leq \theta \leq \theta_1$ ) in the second case; the conditions of diffusion in the domain  $\theta_1 \leq \theta \leq \pi$  are distinguished by extreme complexity because of the presence of a stationary vortex therein [4]. However, the mass exchange of the part of the surface with  $\theta > \theta_1$  can be neglected because this part comprises less than 12.5% of the whole surface of the sphere, and because of the insignificant tangential velocities in this domain.

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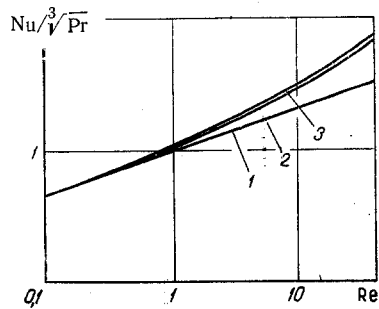


Fig. 1. Critical dependences governing the mass-exchange coefficient: 1) from (1); 2) from (7); (8), 3) from (9).

The solution of the problem under consideration is performed as follows.

1. The tangential velocity  $v_\theta$  and its value within the diffusion layer, i. e., for  $r = a + y$ ,  $y \ll a$ , are determined from (3).

2. The concentration distribution in the diffusion layer is represented by the relationship

$$\frac{c - c_1}{c_s - c_1} = 1 - \sin \frac{\pi}{2} \frac{y}{\delta},$$

which satisfies the boundary conditions

$$c(0) = c_s; \quad c(\delta) = c_1; \quad \left( \frac{\partial c}{\partial y} \right)_{y=\delta} = 0; \quad \left( \frac{\partial^2 c}{\partial y^2} \right)_{y=0} = 0.$$

3. Integration in conformity with (2) and subsequent solution of the differential equation determine the diffusion-layer thickness, and its distribution over the surface

$$\delta = \sqrt[3]{x} \frac{a}{\sqrt[3]{Ua/D}}, \quad (5)$$

$$x = \frac{16.59 \int_0^\theta \sin^2 \theta \left[ \left( 1 + \frac{3}{16} \text{Re} \right) + \frac{\text{Re}}{4} \cos \theta \right]^{1/2} d\theta}{\sin^2 \theta \left[ \left( 1 + \frac{3}{16} \text{Re} \right) + \frac{\text{Re}}{4} \cos \theta \right]^{3/2}}. \quad (6)$$

4. The flux of material from the sphere surface is

$$I = -D \int_0^{\theta_1} \left( \frac{dc}{dy} \right)_{y=0} 2\pi a \sin \theta d\theta$$

and it determines the magnitude of the mass-exchange coefficient

$$I = k(c_s - c_1) 2\pi a^2 (1 - \cos \theta_1).$$

The final result is

$$\begin{aligned} & \text{Re} \leq 16; \quad m \geq 1; \\ \frac{\text{Nu}}{\sqrt[3]{\text{Pr}}} &= 0.304 \sqrt[3]{\text{Re} \left( 1 + \frac{3}{16} \text{Re} \right)} \sqrt[3]{\frac{1+m}{m}} \left[ m(1-m) F \left( \sqrt{\frac{2}{1+m}} \right) \right. \\ & \left. + (3+m^2) E \left( \sqrt{\frac{2}{1+m}} \right) \right]^{2/3}; \quad (7) \end{aligned}$$

$$\begin{aligned} & \text{Re} \geq 16; \quad m \leq 1; \\ \frac{\text{Nu}}{\sqrt[3]{\text{Pr}}} &= 0.412 \sqrt[3]{\text{Re} \left( 1 + \frac{3}{16} \text{Re} \right)} \frac{1}{\sqrt[3]{m(1+m)}} \left[ (m-3) \left( 1 \right. \right. \\ & \left. \left. - m \right) F \left( \sqrt{\frac{1+m}{2}} \right) + (6 + 2m^2) E \left( \sqrt{\frac{1+m}{2}} \right) \right]^{2/3}. \quad (8) \end{aligned}$$

Expanding the elliptic integrals in (7) in series and retaining three terms in each, we obtain result more convenient for use

$$\frac{\text{Nu}}{\sqrt[3]{\text{Pr}}} = 0.99 \sqrt[3]{\text{Re}} \frac{\left(1 + \frac{13}{16} \text{Re} + \frac{17}{128} \text{Re}^2\right)^{2/3}}{1 + \frac{7}{16} \text{Re}} \quad (9)$$

The results (1), (7), (8), (9) are presented in Fig. 1 in logarithmic coordinates. The values determined from (9) exceed the exact value (7) by not more than 5%.

#### NOTATION

$a$	sphere radius;
$c_1$	concentration in the main mass of solution;
$c$	concentration within the diffusion layer;
$c_s$	concentration on the sphere surface;
$F$	complete elliptic integral of the first kind;
$k$	mass-exchange coefficient;
$m$	a parameter defined by (4);
$U$	velocity of fluid motion around the sphere;
$v_\theta$	tangential velocity component;
$x$	a parameter (see (5) and (6));
$y$	distance from the sphere surface in a radial direction $0 \leq y \leq \delta$ ;
$D$	diffusion coefficient;
$r$	radius-vector of a point outside the sphere;
$I$	flux of material from the sphere surface (quantity of material lost by the sphere in unit time);
$\delta$	diffusion-layer thickness;
$\theta$	angular distance from the forward stagnation point;
$\theta_1$	the same at points of vertical-zone formation;
$\nu$	kinematic viscosity;
$\psi$	stream function;
$\text{Nu} = k \cdot 2a/D$	diffusion Nusselt number;
$\text{Pr} = \nu/D$	diffusion Prandtl number;
$\text{Re} = U \cdot 2a/\nu$	Reynolds number;
$E$	complete elliptic integral of the second kind.

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